

Superfluids Under Rotation
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TWISTED VORTEX STATE

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Content

1) Does the twisted state exist at all?

Hydrodynamic theory

- uniform twist
- linear theory of nonuniform twist

2) Generation of twisted vortex states

Numerical simulations

3) Observation in superfluid ^3He



sketch of twisted
vortex lines

Twisted vortex states in classical fluids

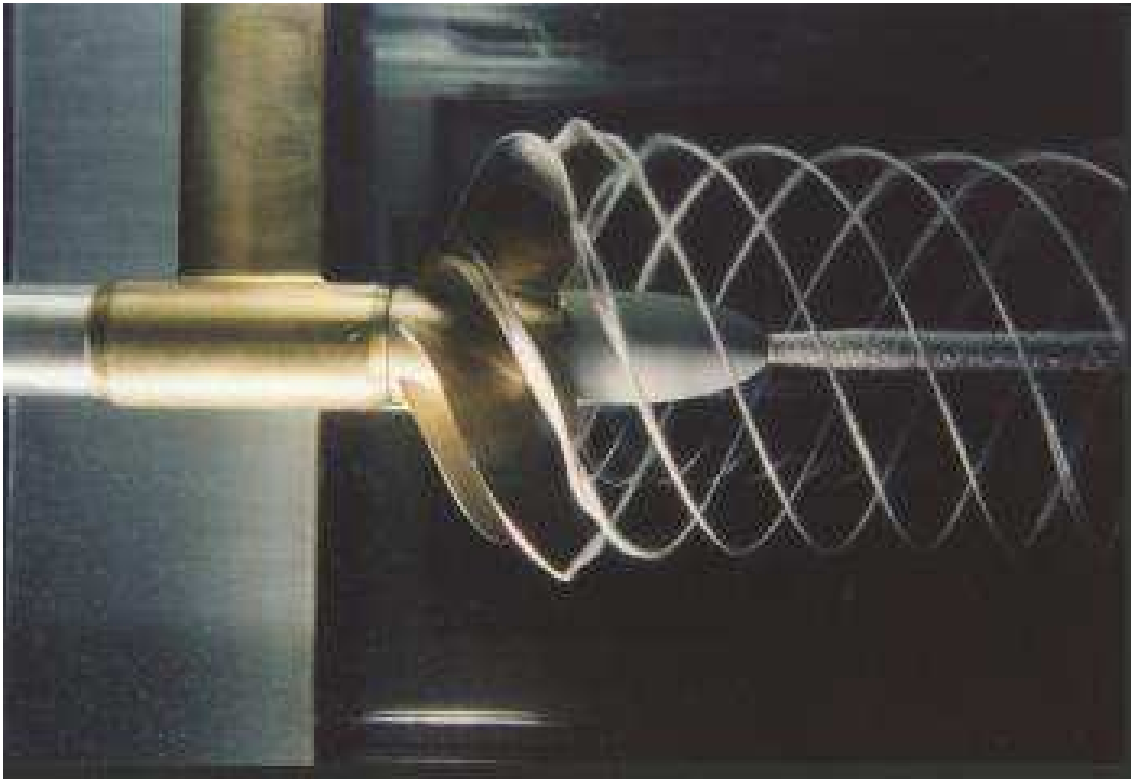
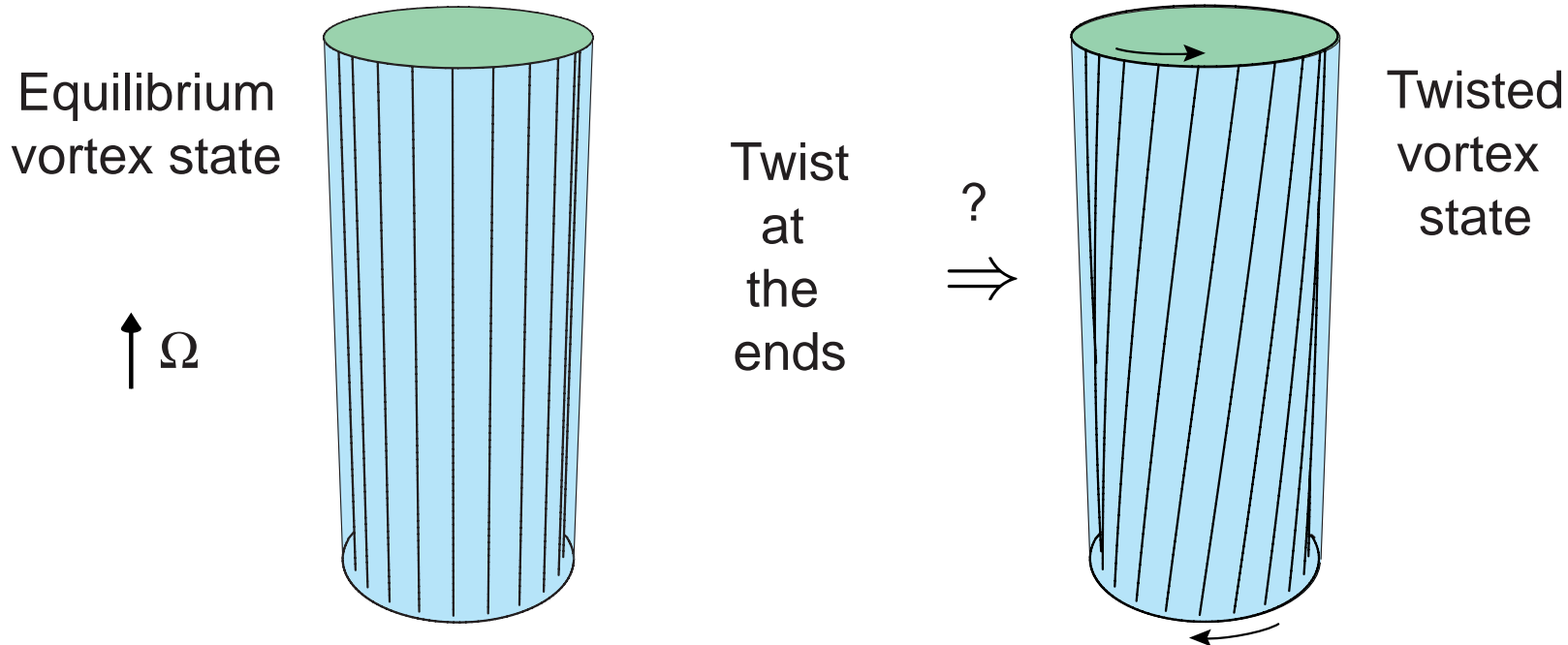


Figure: <http://www.amc.edu.au/research/areas/cavitation/projects/>

Stability of a polygon of helical vortices (Okulov 2004)

Rotating superfluid



Previous literature: Hall (1958), Andronikashvili et al (1961), Glaberson et al (1974), Sonin (1987), Donnelly (1991)...

→ no mention of twisted vortices!

Taylor-Proudman theorem: "Any slow motion in rotating fluid is columnar"

1) Does the twisted state exist?

Hydrodynamic equations

Superfluid velocity v_s

$$\nabla \times v_s = 0 \quad \text{except at vortex lines}$$

$$\nabla \cdot v_s = 0$$

\Rightarrow Vortex lines fully determine $v_s(\mathbf{r}, t)$.

Line velocity v_L

$$v_L = v_s$$

Add mutual friction

$$v_L = v_s + \alpha \hat{l} \times (v_n - v_s) - \alpha' \hat{l} \times [\hat{l} \times (v_n - v_s)] \quad (1)$$

Continuum model of vorticity

(Hall and Vinen 1956)

$$\mathbf{v}_S = \langle \mathbf{v}_S^{\text{local}} \rangle$$

$$\nabla \times \mathbf{v}_S = \boldsymbol{\omega}$$

$$\nabla \cdot \mathbf{v}_S = 0$$

Line velocity \mathbf{v}_L

$$\mathbf{v}_L = \tilde{\mathbf{v}}_S + \alpha \hat{\mathbf{l}} \times (\tilde{\mathbf{v}}_n - \mathbf{v}_S) - \alpha' \hat{\mathbf{l}} \times [\hat{\mathbf{l}} \times (\mathbf{v}_n - \tilde{\mathbf{v}}_S)] \quad (2)$$

where $\tilde{\mathbf{v}}_S = \mathbf{v}_S + \nu \nabla \times \hat{\boldsymbol{\omega}}$, $\nu = (\kappa/4\pi) \ln(b/a)$.

Alternatively, one can use equation of motion for \mathbf{v}_S :

$$\frac{\partial \mathbf{v}_S}{\partial t} = \mathbf{v}_S \times \boldsymbol{\omega} + \nu (\boldsymbol{\omega} \cdot \nabla) \hat{\boldsymbol{\omega}} + \nabla \phi$$

Uniformly twisted vortex state

Most symmetric state [cylindrical coordinates (r, ϕ, z)]

$$\mathbf{v}_s = v_\phi(r)\hat{\phi} + v_z(r)\hat{z},$$

\Rightarrow vorticity

$$\boldsymbol{\omega} = \frac{1}{2}\nabla \times \mathbf{v}_s = \frac{1}{2}\left[-v'_z\hat{\phi} + \left(\frac{v_\phi}{r} + v'_\phi\right)\hat{z}\right]$$

Calculate vortex line velocity from (2). For a stationary state the radial velocity must vanish. This implies

$$(\Omega r - v_\phi)\left(\frac{v_\phi}{r} + \frac{dv_\phi}{dr}\right) - v_z\frac{dv_z}{dr} + \frac{\nu}{|\boldsymbol{\omega}|r}\left(\frac{dv_z}{dr}\right)^2 = 0.$$

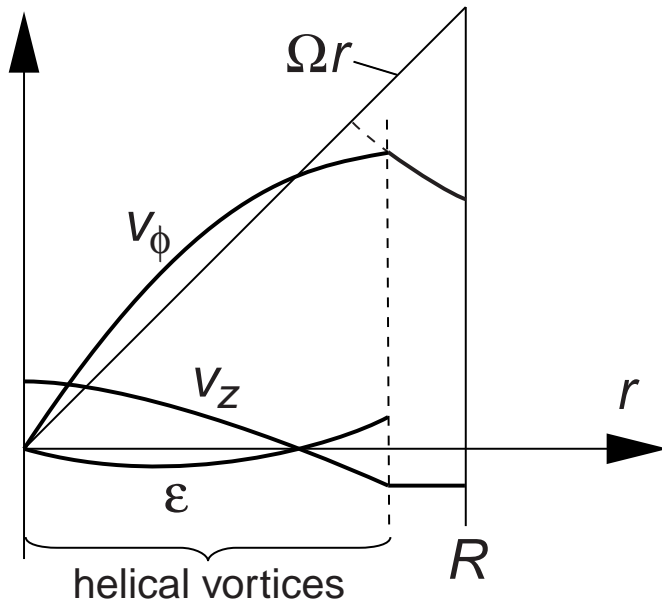
This implies that the helical vortices rotate together with the normal fluid, $\mathbf{v}_L = \mathbf{v}_n = \boldsymbol{\Omega} \times \mathbf{r}$.

\Rightarrow There exists a family of stationary, uniformly twisted states.

In a finite cylinder the total axial current must vanish,

$$\int_0^R dr r v_z = 0. \quad (3)$$

The functions $v_z(r)$, $v_\phi(r)$ and the radial displacement of the vortices compared to equilibrium state, $\epsilon(r)$, are sketched in the figure.



The simplest case is helical vortices with a wave vector $Q(r) = \text{constant}$. This has

$$\begin{aligned} v_\phi(r) &= \frac{(\Omega + Qv_0)r}{1 + Q^2r^2}, \\ v_z(r) &= \frac{v_0 - Q\Omega r^2}{1 + Q^2r^2}. \end{aligned} \quad (4)$$

Linearized hydrodynamics

Assume general velocity with circular symmetry

$$\mathbf{v}_s = v_r(r, z, t)\hat{\mathbf{r}} + v_\phi(r, z, t)\hat{\boldsymbol{\phi}} + v_z(r, z, t)\hat{\mathbf{z}}$$

Assume small deviation from rotating equilibrium.

⇒ waves of the form

$$v_r = ckJ_1(\beta r) \exp(ikz - i\sigma t)$$

$$v_z = ic\beta J_0(\beta r) \exp(ikz - i\sigma t)$$

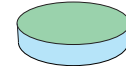
Dispersion relation [Glaberson, Johnson and Ostermeier (1974), Henderson and Barenghi (2004)]

$$\frac{\sigma}{\Omega} = \frac{-i\alpha(\beta^2 + 2k^2\eta_2) \pm i\sqrt{\alpha^2\beta^4 - 4(1 - \alpha')^2k^2(\beta^2 + k^2)\eta_1\eta_2}}{\beta^2 + k^2}$$

where $\eta_1 = 1 + \nu k^2/2\Omega$ and $\eta_2 = 1 + \nu(\beta^2 + k^2)/2\Omega$.

In order to understand the dispersion, we study special cases.

1) $\beta \rightarrow 0$, corresponds to a short cylinder



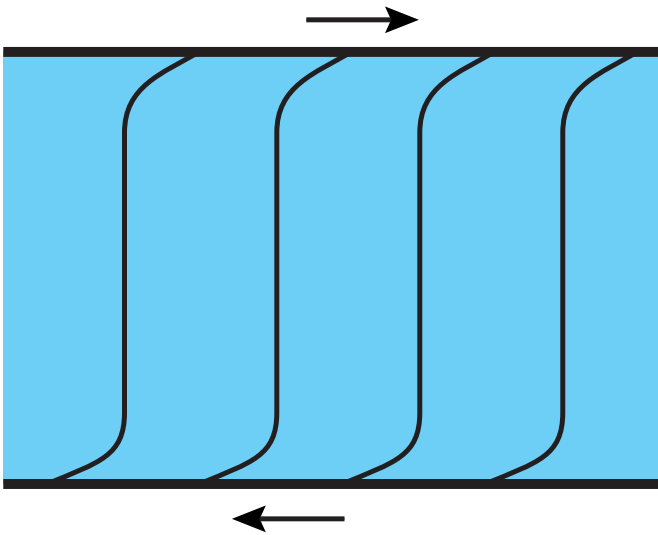
\Rightarrow 2 Kelvin wave modes (Hall 1958)

$$k_{\pm} = i\sqrt{\frac{2\Omega \pm \sigma}{\nu}}$$

and an inertial mode

$$k_i = 0$$

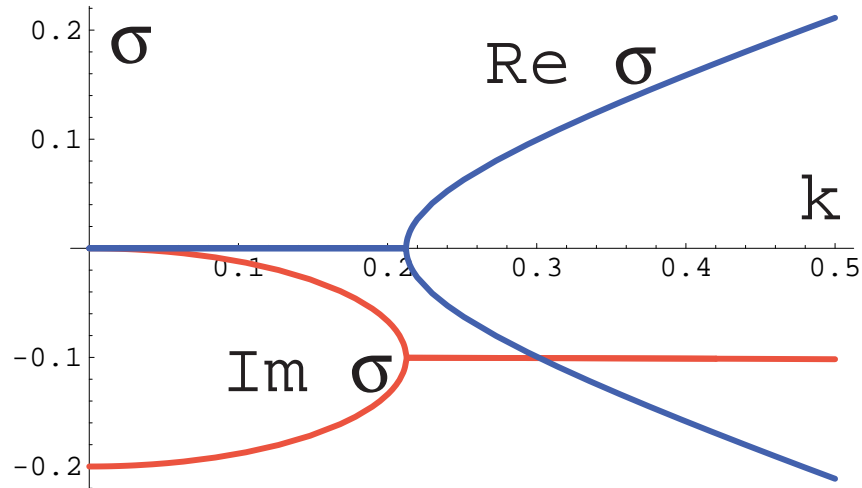
At low frequency ($\sigma \ll \Omega$) these give just the columnar motion because Kelvin waves are evanescent. No twisted state.



2) $k \rightarrow 0$, corresponds to a long cylinder



\Rightarrow 2 modes



The point $k = \sigma = 0$ corresponds to uniform twist!

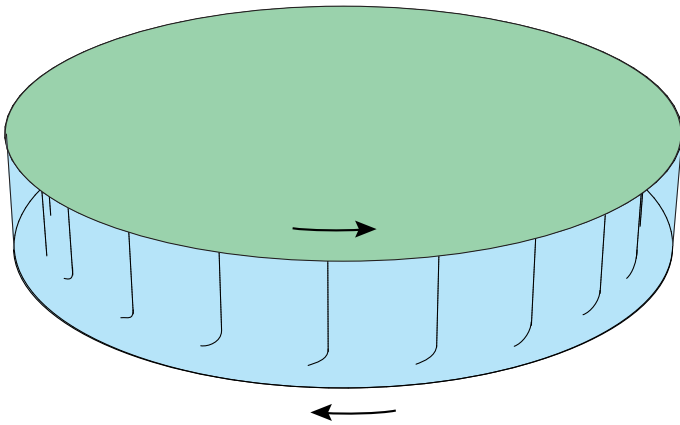
At finite k the twist obeys diffusion equation

$$\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial z^2}, \quad D = \frac{1}{d} \left(\frac{2\Omega}{\beta^2} + \nu \right) \quad (5)$$

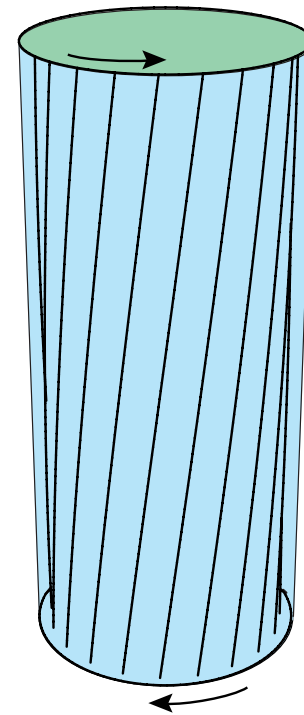
where $f(z, t) = v_r$ or v_z .

Summary of two opposite limits

Parallel plates
- columnar vortices



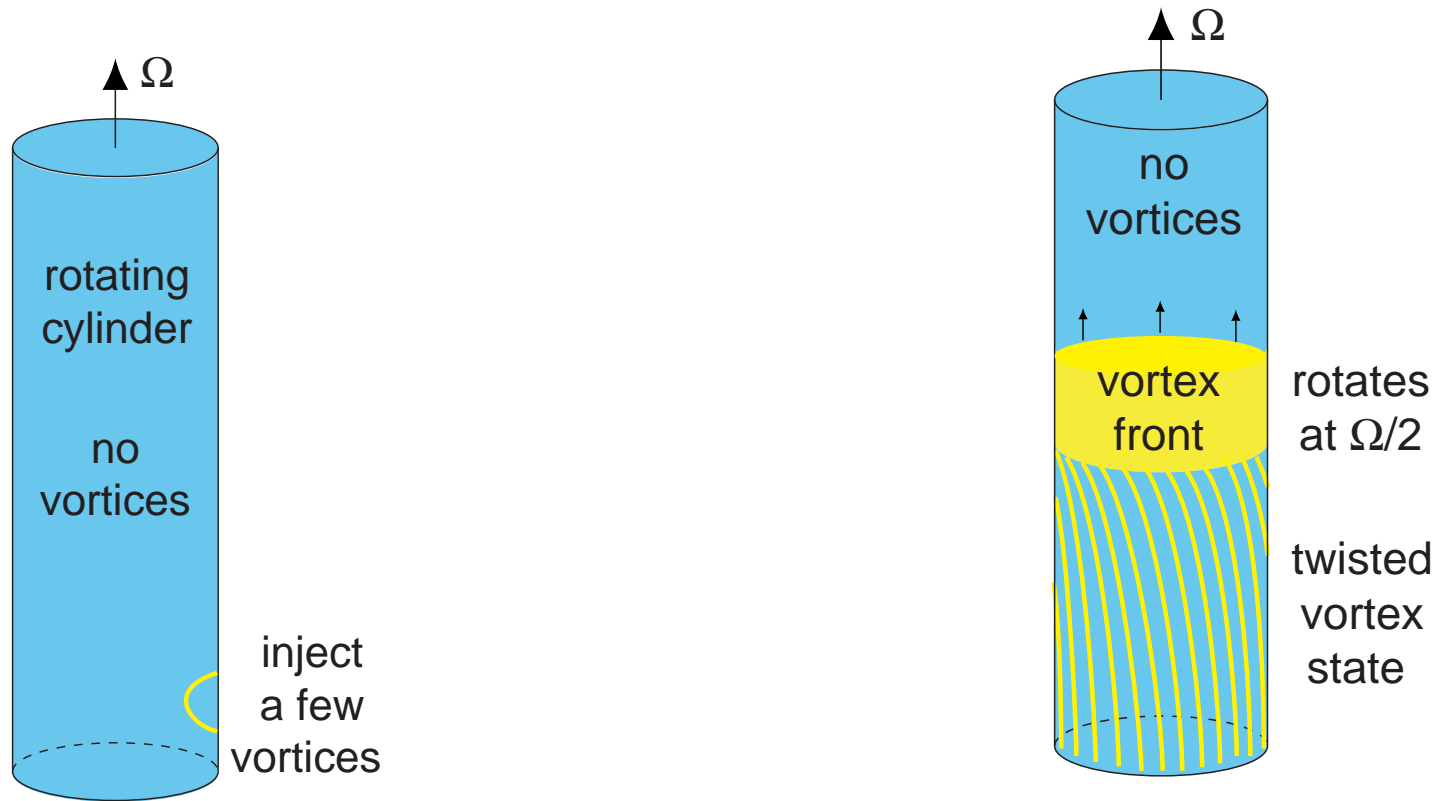
Long cylinder
- twisted vortices



2) Generation of twisted vortex states

- superfluid in a cylinder
- cylinder rotating at $\Omega > \Omega_c$,
but
no vortices in the initial state
- generate vortices at one place

- vortices propagate along the cylinder and
- vortex ends rotate around the cylinder axis



Why vortex ends rotate?

Normal component rotates at $\mathbf{v}_n = \boldsymbol{\Omega} \times \mathbf{r}$.

Superfluid component: vortex lines move with the average superfluid velocity

1) vortex state: $\mathbf{v}_s \approx \boldsymbol{\Omega} \times \mathbf{r}$

\Rightarrow vortex lattice rotates at angular velocity $\boldsymbol{\Omega}$

2) no vortices: $\mathbf{v}_s = 0$

3) vortex front

average superfluid angular velocity $\boldsymbol{\Omega}/2 \Rightarrow$ vortex ends rotate at angular velocity $\boldsymbol{\Omega}/2$

\Rightarrow propagating vortex ends lag behind

Numerical simulation

Vortex line velocity (2)

$$\mathbf{v}_L = \mathbf{v}_s + \alpha' \hat{\mathbf{l}} \times [(\mathbf{v}_n - \mathbf{v}_s) \times \hat{\mathbf{l}}] + \alpha \hat{\mathbf{l}} \times (\mathbf{v}_n - \mathbf{v}_s).$$

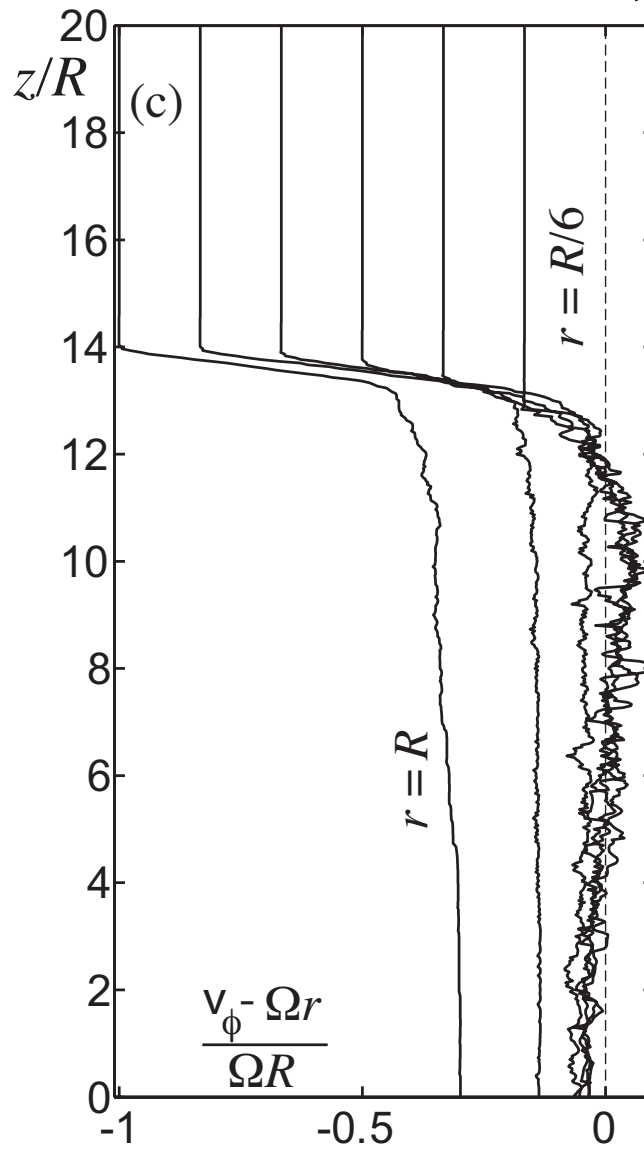
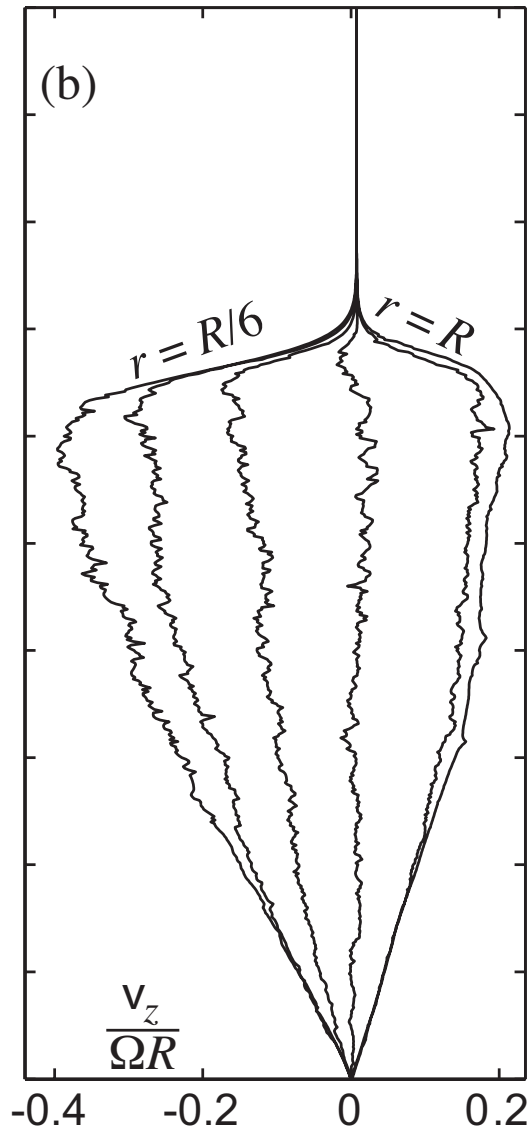
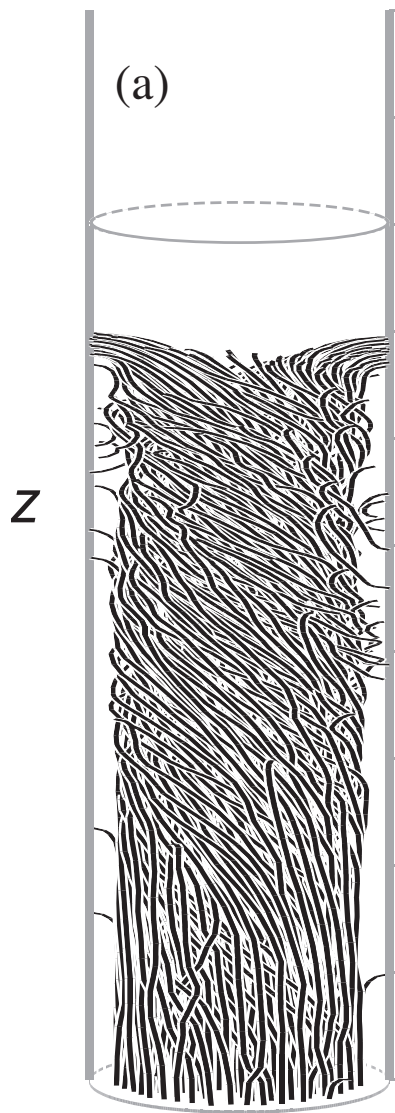
\mathbf{v}_s is calculated from Biot-Savart integral. (Risto Hänninen)

The front and the twisted state is confirmed by numerical calculation

movie

Axial velocity

Azimuthal velocity



movie

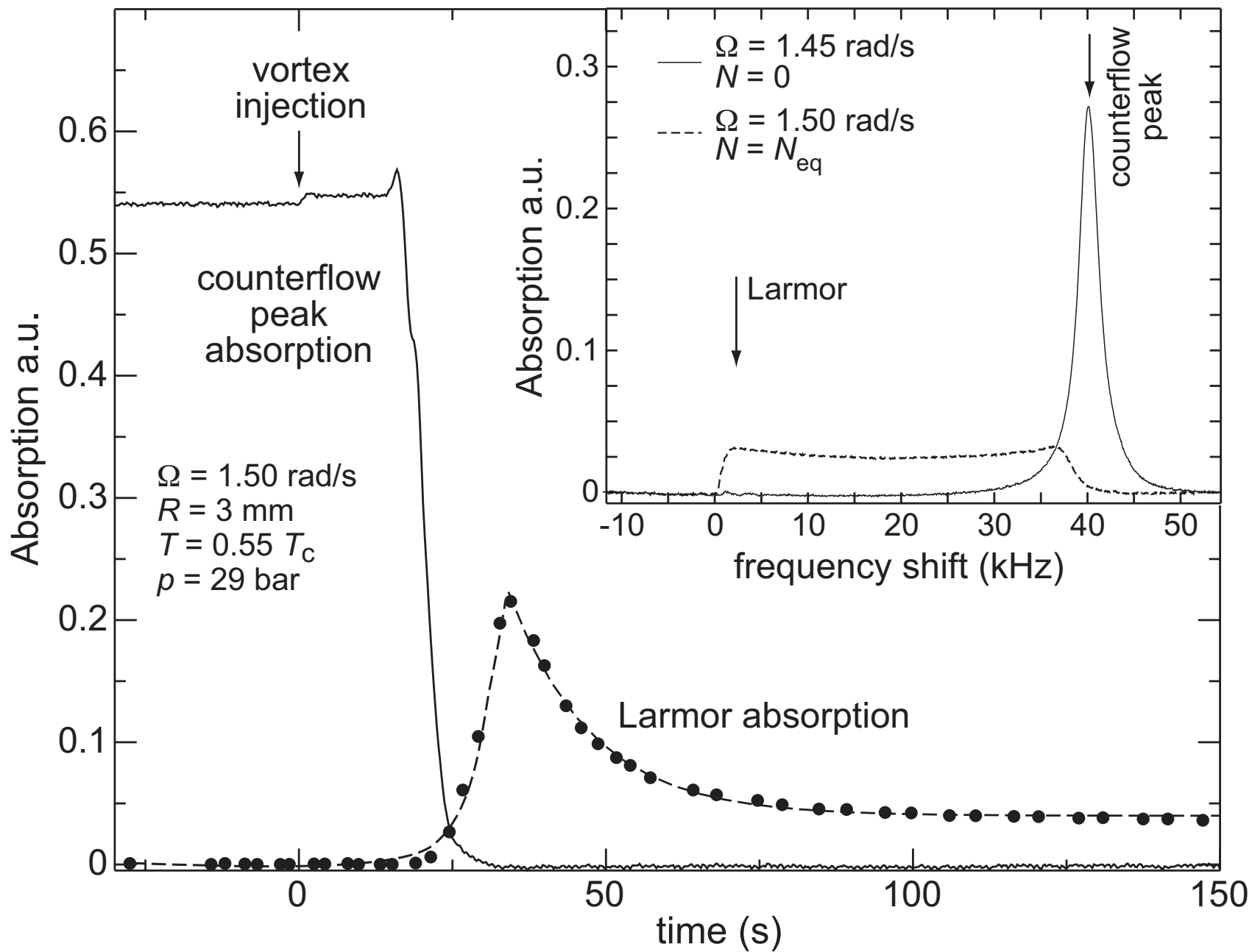
Main observations

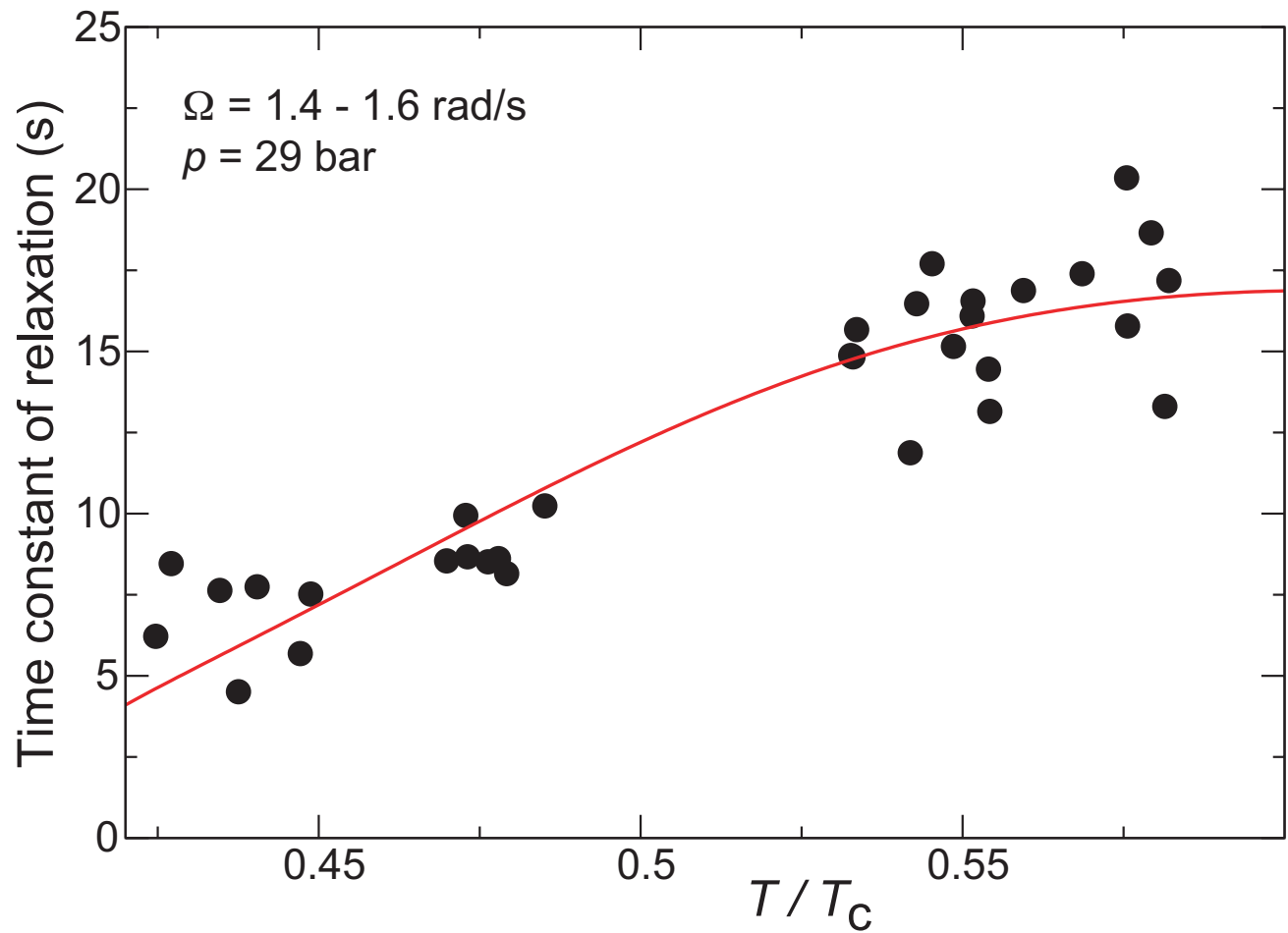
- the twisted state has axial current.
- individual vortices become unstable to generate Kelvin waves at large axial current
- the vortices glide at the bottom plate
⇒ relaxation of the twist
- the relaxation is determined by the diffusion equation.

3) Experiment in superfluid $^3\text{He-B}$

Vortex state was generated as discussed above.

The axial velocity v_z affects the texture, which is seen by NMR.





Diffusion constant

$$D = \frac{1}{d} \left(\frac{2\Omega}{\beta^2} + \nu \right) \propto \frac{1}{\text{mutual friction constant}} \quad (6)$$

Conclusions

Twisted vortex state is a possible state in long rotating cylinders.

The twisted state can be generated by vortex injection.

The twisted state has been seen in superfluid $^3\text{He-B}$.

Eltsov et al, Phys. Rev. Lett. 96, 215302 (2006)